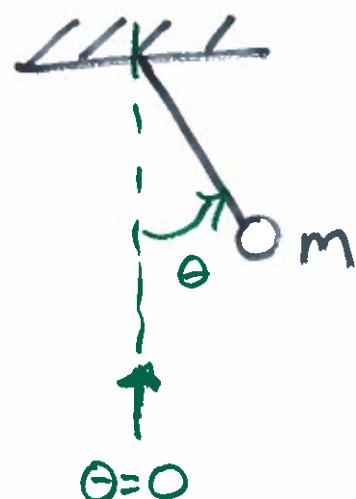


Lecture 3 Wednesday, September 28, 2016
Linearized EoM of a simple pendulum



equilibrium point

inertial force
restoring force

1 DoF

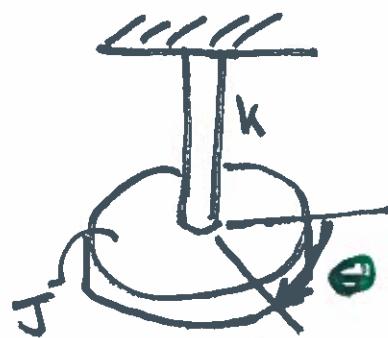
second order linear ODE

$$\ddot{\theta} + \frac{g}{l} \theta = 0 \quad (1)$$

equivalent to:



$$\ddot{x} + \frac{k}{m} x = 0$$



$$\ddot{\theta} + \frac{k}{J} \theta = 0$$

Solution to ODE

Sol of (1) $\theta(t) = A \sin(\omega_n t + \phi)$

?

A: amplitude

ω_n : natural freq

ϕ : phase shift

$$\dot{\theta}(t) = \omega_n A \cos(\omega_n t + \phi)$$

$$\ddot{\theta}(t) = -\omega_n^2 A \sin(\omega_n t + \phi)$$

$$\ddot{\theta} + \frac{g}{l} \theta = 0$$

$$-\omega_n^2 A \sin(\omega_n t + \phi) + \frac{g}{l} A \sin(\omega_n t + \phi) = 0$$

$$\omega_n^2 = \frac{g}{l} \Rightarrow \boxed{\omega_n = \sqrt{\frac{g}{l}}} \quad \text{natural frequency}$$

A and ϕ depends initial conditions.

$$1) \theta_0 = \theta(0) = A \sin(\omega_n \cdot 0 + \phi) = A \sin \phi$$

$$2) \omega_0 = \dot{\theta}(0) = \omega_n A \cos(\omega_n \cdot 0 + \phi) = \omega_n A \cos \phi$$

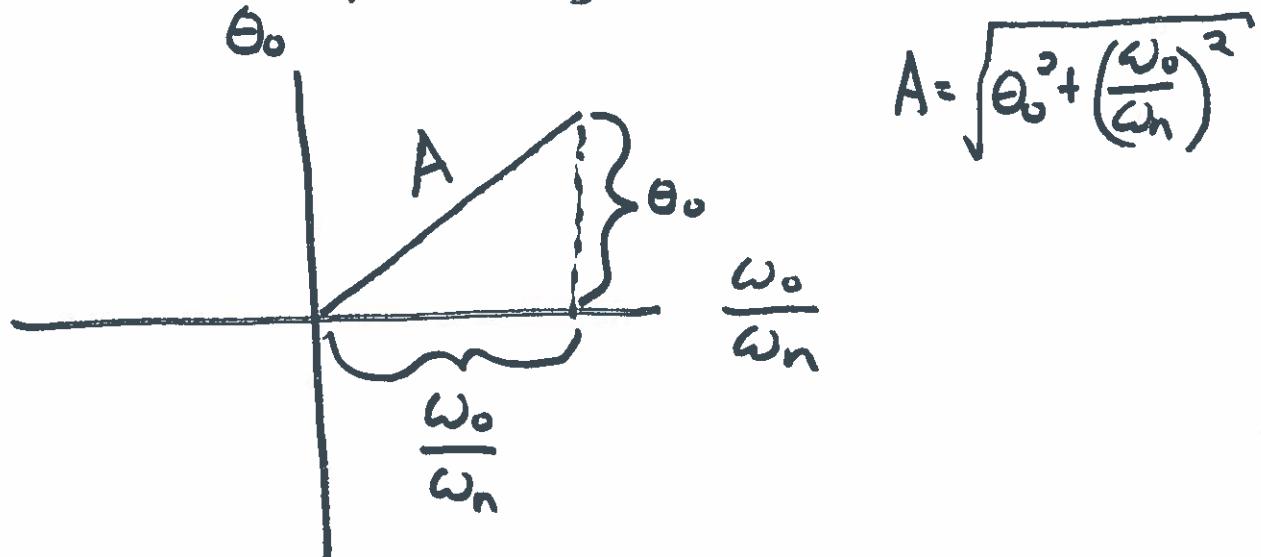
Solve for A and ϕ :

$$A = \sqrt{\frac{\omega_n^2 \theta_0^2 + \omega_0^2}{\omega_n}} \quad \phi = \arctan\left(\frac{\omega_n \theta_0}{\omega_0}\right)$$

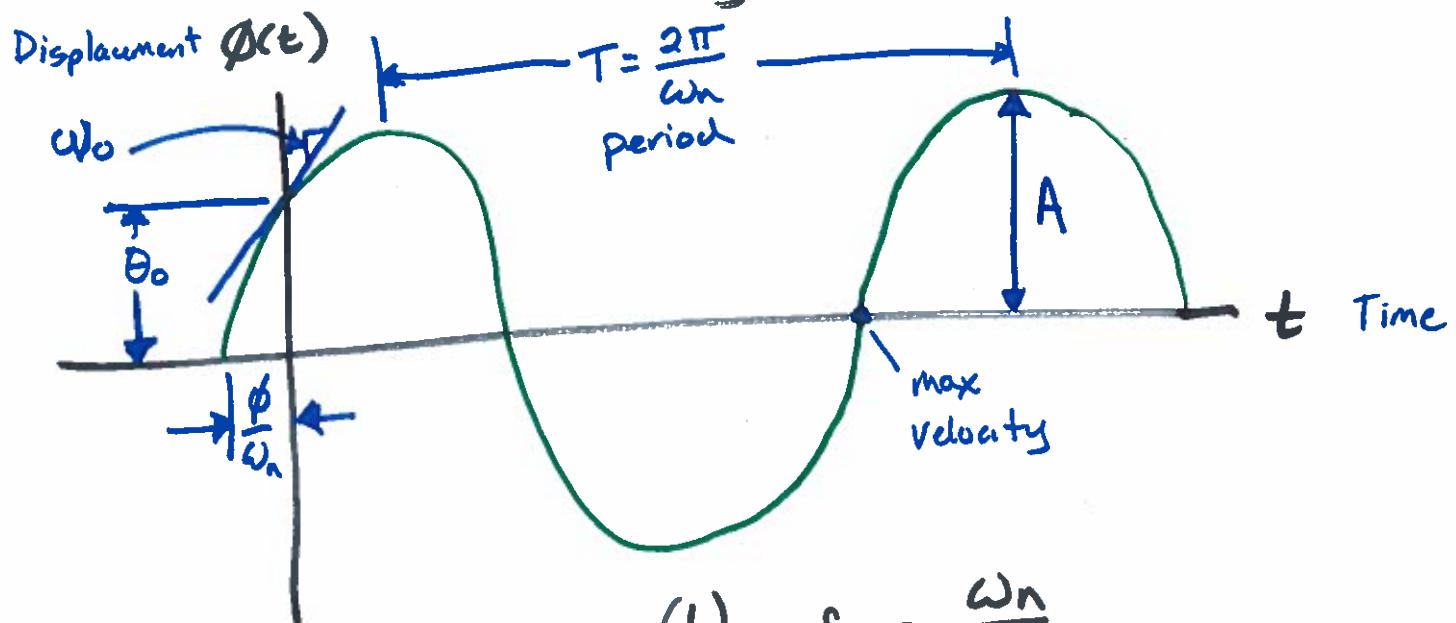
Full equation

$$\theta(t) = \sqrt{\frac{\omega_n^2 \theta_0^2 + \omega_0^2}{\omega_n}} \sin\left[\omega_n t + \arctan\left(\frac{\omega_n \theta_0}{\omega_0}\right)\right]$$

Geometric relationships among solution unknowns.



Beware of the quadrant this falls in when computing the arctan !!

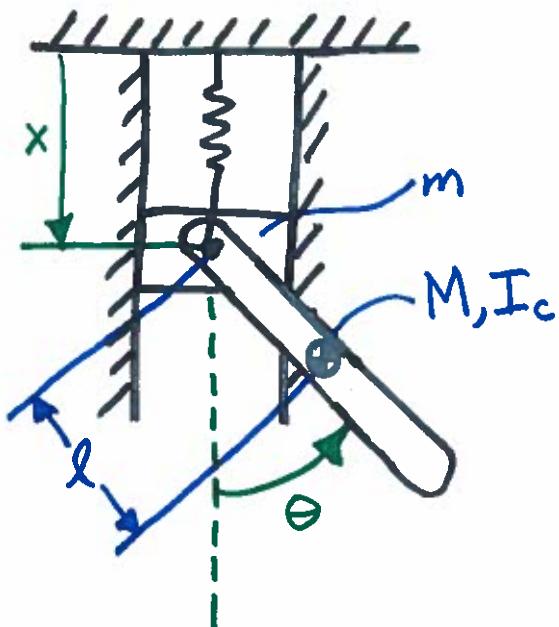


f_n : frequency in Hertz ($\frac{1}{s}$) $f_n = \frac{\omega_n}{2\pi}$

T: period in seconds $T = \frac{1}{f_n}$

I.C.s contain info about initial energy in system

Corrected wrt lecture

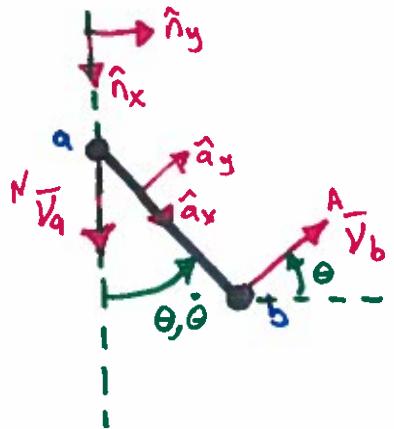


Example: Compound pendulum with base on vertical spring

Goal: Find the equations of motion for this model using Lagrange's method.

Model has 2 DoF with generalized coordinates x and θ .

Step 1: Form an expression for the Kinetic energy



$${}^N\bar{v}_b = {}^N\bar{v}_a + {}^A\bar{v}_b$$

$${}^A\bar{v}_b = {}^N\omega^A \times \bar{r}_{b/a}$$

N: inertial ref. frame
A: ref. frame attached to bar ab

Note the corrected magnitude of the velocity at point b!

$$|{}^N\bar{v}_b| = |{}^N\bar{v}_a + {}^A\bar{v}_b|$$

$$= \sqrt{(\dot{x} - \dot{\theta}l \sin \theta)^2 + (\dot{\theta}l \cos \theta)^2}$$

$$|{}^N\bar{v}_b| = \sqrt{\dot{x}^2 + l^2\dot{\theta}^2 - 2\dot{x}l\dot{\theta}\sin\theta}$$

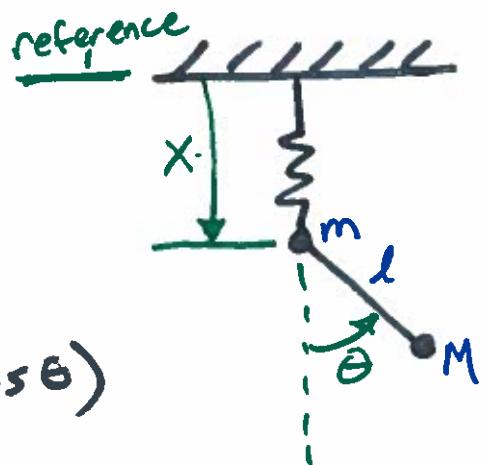
$$T = \frac{1}{2}m|\bar{v}_a|^2 + \frac{1}{2}M|\bar{v}_b|^2 + \frac{1}{2}I_c\dot{\theta}^2$$

$$T = \underbrace{\frac{m}{2}\dot{x}^2 + \frac{M}{2}(\dot{x}^2 + l^2\dot{\theta}^2 - 2\dot{x}l\dot{\theta}\sin\theta)}_{\text{linear K.E.}} + \underbrace{\frac{I_c}{2}\dot{\theta}^2}_{\text{rotational Kinetic energy}}$$

linear K.E.

rotational
kinetic
energy

$$U = U_{\text{spring}} + \underbrace{U_{mg} + U_{Mg}}_{\text{P.E. due to gravity}}$$



$$U = \frac{1}{2} k x^2 - mgx - Mg(x + l \cos \theta)$$

Step 2: Form Lagrangian

$$L = T - U$$

$$L = \frac{m}{2} \dot{x}^2 + \frac{M}{2} (\dot{x}^2 + l^2 \dot{\theta}^2 - 2\dot{x}l\dot{\theta}\sin\theta) + \frac{I_c}{2} \dot{\theta}^2 - \frac{k}{2} x^2 + mgx + Mg(x + l\cos\theta)$$

Step 3: Form Lagrange's Equations

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}_i} \right) - \frac{\partial L}{\partial x_i} = 0$$

$$x_1 = x, x_2 = \theta$$

two D.o.F

two G.C.s

two second order
coupled non-linear
ODEs

① X

$$\frac{d}{dt} (m\ddot{x} + M\ddot{x} - Ml\dot{\theta}\sin\theta) + Kx - mg - Mg = 0$$

$$m\ddot{x} + M\ddot{x} - M(l\ddot{\theta}\sin\theta + l\dot{\theta}^2\cos\theta) + Kx - mg - Mg = 0$$

(5)

② θ

$$\frac{d}{dt} (Ml^2\ddot{\theta} - M\dot{x}lsin\theta + I_c\ddot{\theta}) + M\dot{x}l\dot{\theta}cos\theta + Mgl sin\theta = 0$$

$$Ml^2\ddot{\theta} - M\ddot{x}lsin\theta - M\dot{x}\dot{\theta}cos\theta + I_c\ddot{\theta} + M\dot{x}l\dot{\theta}cos\theta + Mgl sin\theta = 0$$

$$Ml^2\ddot{\theta} - M\ddot{x}lsin\theta + I_c\ddot{\theta} + Mgl sin\theta = 0$$

We now have two coupled second order non-linear ordinary differential equations. With some simplification we have:

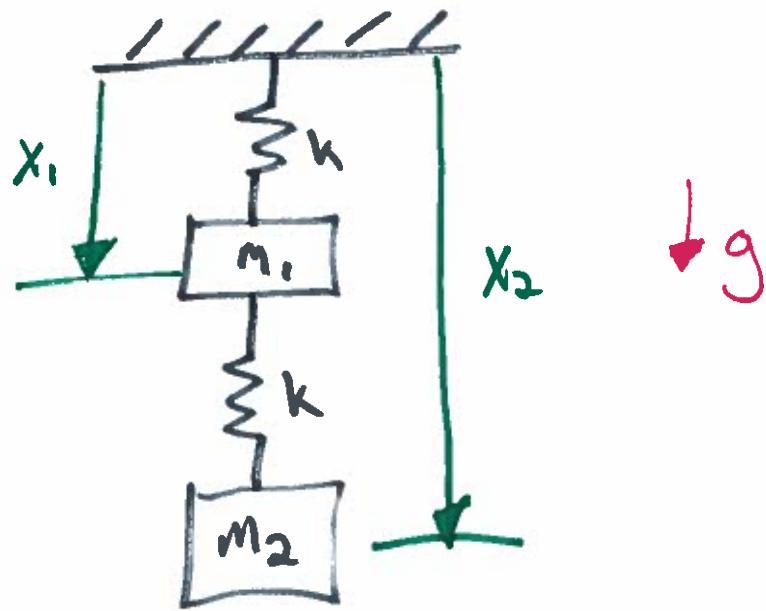
$$(m+M)\ddot{x} - \underbrace{Ml(\ddot{\theta}sin\theta + \dot{\theta}^2cos\theta)}_{\begin{array}{l} \text{force required} \\ \text{to linearly accel.} \\ \text{both masses} \end{array}} + \underbrace{Kx}_{\begin{array}{l} \text{force req.} \\ \text{to centripetally} \\ \text{accel bar} \end{array}} - \underbrace{(m+M)g}_{\begin{array}{l} \text{force due to} \\ \text{gravity} \end{array}} = 0$$

$$(ml^2 + I_c)\ddot{\theta} - \underbrace{M\dot{x}lsin\theta}_{\begin{array}{l} \text{torque req.} \\ \text{to angularly} \\ \text{accel bar} \end{array}} + \underbrace{Mgl sin\theta}_{\begin{array}{l} \text{torque required for bar} \\ \text{to overcome} \\ \text{gravity} \end{array}} = 0$$

⑥

Exercise

Write out the potential energy, U , term for the following system:



Do not assume springs are at equilibrium

Both systems are conservative \Rightarrow no loss of energy
simply transforms kinetic to potential and vice versa.

So

$T+U$ is constant wrt to time
total energy

(7)

$$U = U_{s_1} + U_{s_2} + U_{g_1} + U_{g_2}$$

s: spring
g: gravity

$$U = \frac{K}{2}x_1^2 + \frac{K}{2}(x_2 - x_1)^2 - m_1 g x_1 - m_2 g x_2$$

What is the equilibrium point?

Spring 1 supports weight of m_1 and m_2 so:

$$x_1^{eq} = \frac{(m_1 + m_2)g}{K}$$

Spring 2 only supports weight of m_2 so:

$$x_2^{eq} = x_1^{eq} + \frac{m_2 g}{K} = \frac{(m_1 + 2m_2)g}{K}$$

The potential energy stored at equilibrium can be found by substituting the equilibrium point into U .